



half-metals – high spin-polarization and broad variability of their magnetization combined with a strong magnetocrystalline anisotropy. Fig.1 shows how different chemical substitutions of Mn in  $Mn_3Ga$  by other  $3d$  transition elements lead to the localization of the spin-down electronic states in the vicinity of the Fermi energy. By the first-principles calculations of the spin-resolved residual conductivities based on the Kubo-Greenwood formalism, we have shown that the spin-polarization  $P_z$  (along the  $c$ -axis of a tetragonal unit cell) might be dramatically increased by using already a very small substitution rates. In the metallic regime  $P_z$  depends much stronger on the ratio of the spin-resolved electron mobilities rather than on spin-resolved concentrations of conducting electrons. Our latest studies on other Heusler systems, for example  $Mn_{3-x}Pt_xGa$  series, indicate that the spin-selective localization is rather widely spread phenomenon, which occurs in a large number of tetragonal and cubic Mn-based Heusler alloys.

**Noncollinear magnets.** Noncollinear distribution of the magnetization gains nowadays a special focus initiated by novel principles of the MRAM technologies. Especially interesting are the so-called magnetic Skyrmion structures: long-range vortices formed by the interplay of several factors, among which the spin-orbit coupling and the absence of the inversion symmetry are the most necessary ingredients. Recently, we have computed and analyzed these ingredients for the group of tetragonal Mn-based alloys [4], which suggests the Skyrmion order in  $Mn_2RhSn$  Heusler material (Fig.2):

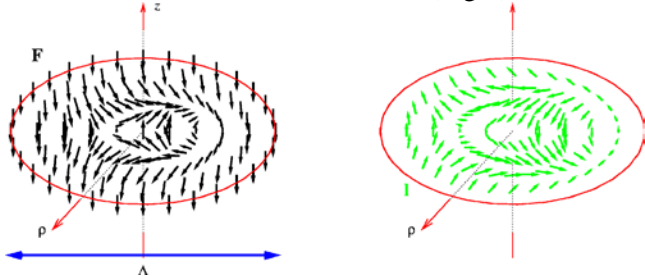


Fig. 2: Schematic view of the magnetic skyrmion phase in  $Mn_2RhSn$ . The  $z$ -axis is oriented along the tetragonal  $c$ -axis. Right: ferromagnetic mode; left: antiferromagnetic mode. The chiral twisting length estimated by micromagnetic model is  $A \sim 130$  nm.

Antiparallel alignment of the magnetic moments in two different Mn sublattices is typically encountered in most of the  $Mn_2YZ$  compounds, but not all of them exhibit noncollinearity. The most significant exchange coupling between the nearest  $Mn_I$  and  $Mn_{II}$  atoms in different Wyckoff sites is characterized by a large exchange constant ( $J_1 \sim -20$  meV) that leads to a typical collinear ferrimagnetic state. However, we have identified several systematic cases (tetragonal  $Mn_2RhSn$ ,  $Mn_2PtIn$  and  $Mn_2IrIn$ ), in which the collinear ferrimagnetic order cannot explain the

measured saturated magnetization. In this systems we figured out that their collinear order can be perturbed by the next important interaction  $j$  between the next-nearest planes, e.g., between  $Mn_{II}$ - $Y$  planes (Fig.3).

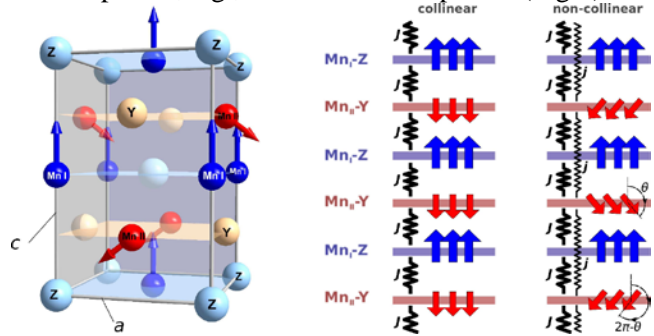


Fig 3.  $Mn_2RhSn$  unit cell: two types of Mn atoms (red and blue); arrows indicate their magnetic moments. Classical exchange picture explains the short-range magnetic canting: additional antiparallel coupling ( $j$ ) competes with the strongest nearest neighbor antiparallel coupling ( $J$ ).

This interaction is antiparallel due to its indirect origin realized through the main-group element Z (super-exchange). Since  $j$  tends to rotate the moments of the nearest  $Mn_{II}$ - $Y$  planes antiparallel to each other, it competes with the strong antiparallel exchange  $J$  and may then result in a nontrivial canting angle ( $\theta \neq 0^\circ; 180^\circ$ ). The relevant  $\theta$ -dependent part of the Heisenberg Hamiltonian contains only two types of antiparallel interactions  $H(\theta) = -J \cdot \cos\theta - \frac{1}{2}j \cdot \cos 2(\pi - \theta)$ , the first term being the coupling of the nearest planes ( $Mn_I$ - $Z$  with  $Mn_{II}$ - $Y$ ) and the second – of the next-nearest ( $Mn_{II}$ - $Y$ ) ones. The extrema of the  $H(\theta)$  function,  $\theta_{1,2} = 180^\circ \pm \arccos(J/2j)$ , give a noncollinear solution for  $j > \frac{1}{2}J$ , i.e. the canting occurs only if the next-nearest exchange  $j$  is sufficiently strong. Parameterized by the  $ab$ -intio  $J$  and  $j$  exchange constants, this model reasonably describes the  $\theta$ -dependence of the total energy exhibiting two minima at  $\theta_{1,2} = 180^\circ \pm 55^\circ$ , which was also confirmed by neutron diffraction ( $\theta_{1,2} = 180^\circ \pm 58.9^\circ$ ).

## References

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