

## Oxides with peculiar electronic and magnetic properties

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**The Functional Oxides group explores oxides with unusual electronic and magnetic properties that are a manifestation of the delicate balance between charge, spin, orbital, and lattice degrees of freedom. An important result was the discovery of a Verwey-type charge ordering transition in  $\text{Cs}_4\text{O}_6$ , a compound with mixed-valence negatively charged oxygen molecules in its crystal structure.  $\text{Cs}_4\text{O}_6$  shows a structural transition from a cubic charge disordered to tetragonal charge ordered phase with distinct superoxide  $\text{O}_2^-$  and peroxide  $\text{O}_2^{2-}$  anions, which is accompanied by a drastic change in electrical conductivity. A similar transition was found for  $\text{Rb}_4\text{O}_6$ , where spin dimerization occurs at low temperatures, possibly driven by orbital ordering. The new layered oxoruthenates(V) with low-dimensional structural features,  $\beta\text{-Ag}_3\text{RuO}_4$  and  $\text{AgRuO}_3$ , were synthesized by hydrothermal methods. While  $\beta\text{-Ag}_3\text{RuO}_4$  exhibits cluster-like magnetism in the paramagnetic regime and magnetic ordering below 80 K,  $\text{AgRuO}_3$  apparently lacks any long-range magnetic order despite strong antiferromagnetic exchange interactions.**

Oxides with open-shell electronic systems attract a lot of research interest due their various physical properties, which are a consequence of competing interactions with comparable energy scales. The particular balance between kinetic energy gain accompanying electron delocalization, electron correlation, electron-lattice interactions, magnetic exchange interactions, and spin-orbit coupling gives rise to phenomena like charge and orbital ordering, insulator-to-metal transitions, spin-orbit assisted Mott-insulating states, frustrated magnetism, colossal magnetoresistance, multiferroicity, and high-temperature superconductivity. In this report, we summarize our work on two peculiar oxide systems: (i) oxides with anionic molecular units featuring an open-shell  $p$  electron configuration, and (ii) new ruthenates with layered crystal structures.

Recently, commonalities have been pointed out between strongly correlated transition metal compounds and molecular solids with open-shell  $\text{O}_2^-$  anions. A prototype system is the superoxide  $\text{CsO}_2$ , where low-dimensional magnetic order driven by orbital order was established. Our nuclear magnetic resonance (NMR) [1] and electron paramagnetic resonance (EPR) [2] spectroscopy studies revealed  $\text{CsO}_2$  to be a versatile model system for one-dimensional quantum antiferromagnetism; here, the low-temperature ground state could be considered as a spin Tomonaga-Luttinger liquid, while at higher temperatures, phonon-modulated magnetic interactions occur between spins on a square lattice. The alkali sesquioxides,  $\text{A}_4\text{O}_6$  ( $\text{A} = \text{Cs}, \text{Rb}$ ), which contain one diamagnetic peroxide  $\text{O}_2^{2-}$  and two paramagnetic  $\text{O}_2^-$  units per formula unit and thus constitute anionic mixed-valence systems, extend the analogy between  $3d$  and  $2p$  open-shell systems by adding the charge degree of freedom. Our previous

study revealed the existence of a charge-ordered tetragonal variant of  $\text{Rb}_4\text{O}_6$  [3] in addition to the known cubic structure, whereas EPR and NMR studies on  $\text{Cs}_4\text{O}_6$  indicated a structural transition from the high-temperature cubic to a lower-symmetry phase [4]. We now could resolve the detailed nature of the phase transitions in alkali sesquioxides by a unique combination of experimental techniques. It was demonstrated that the sesquioxides feature a Verwey-type charge ordering (CO) transition that can be compared with that of the canonical example,  $\text{Fe}_3\text{O}_4$  (magnetite) [5].

Low-dimensional magnetism in layered crystal structures provides basic insights into the nature of exchange interactions and often gives rise to peculiar magnetic properties due to spin frustration and quantum spin fluctuations. Herein we report the synthesis and magnetic properties of two new layered ruthenates with low-dimensional structural and magnetic features.

### Verwey-type charge ordering transition in alkali sesquioxides $\text{A}_4\text{O}_6$ ( $\text{A} = \text{Cs}, \text{Rb}$ )

The Verwey transition in  $\text{Fe}_3\text{O}_4$ , a structural phase transition concomitant with a jump in electrical conductivity by two orders of magnitude, has been a benchmark for CO phenomena in mixed-valence transition metal materials. The CO pattern in  $\text{Fe}_3\text{O}_4$ , however, is complex and the mechanism of Verwey transition is still not clear. We have demonstrated that the open  $p$ -shell mixed-valence compound  $\text{Cs}_4\text{O}_6$  shows an archetypical Verwey-type CO transition in the anionic sublattice [5]. Using powder neutron diffraction, it was established that  $\text{Cs}_4\text{O}_6$  undergoes a phase change from a cubic structure with a single crystallographic site for the molecular  $\text{O}_2^{x-}$  building

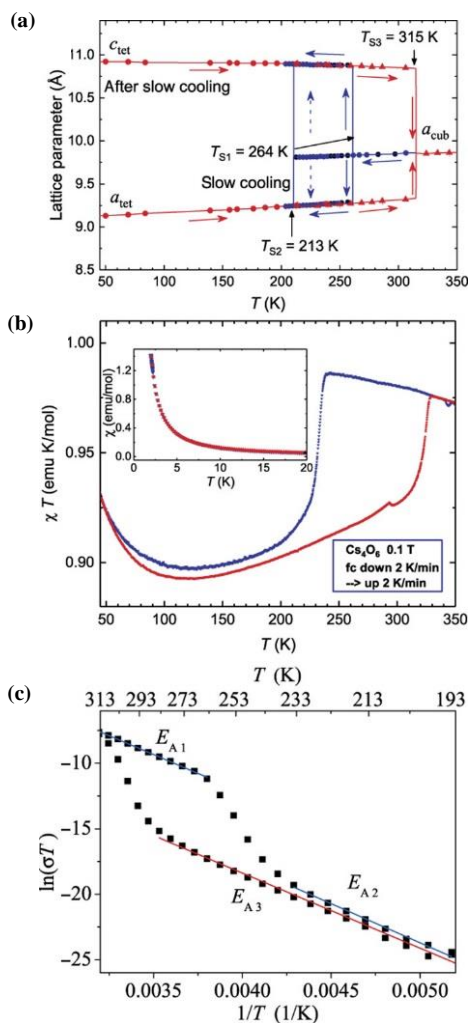


Fig.-1: (from Ref. [5]). Signatures of the Verwey-type CO transition in  $Cs_4O_6$ . (a) Temperature dependence of the lattice parameters:  $a_{cub}$  of the cubic, and  $a_{tet}$  and  $c_{tet}$  of the tetragonal phases. On cooling, the cubic to tetragonal transformation starts at  $T_{S1} = 264$  K and is completed at  $T_{S2} = 213$  K. Upon heating, the transformation back to the cubic phase occurs near  $T_{S3} = 315$  K. (b) The structural transition is reflected in anomalies and a hysteresis in the magnetic susceptibilities  $\chi(T)$ , as apparent in the representation of  $\chi T$  versus  $T$ . The inset shows the low-temperature magnetic susceptibilities verifying the absence of long-range magnetic order down to 1.8 K. (c) Arrhenius plot of the conductivity of  $Cs_4O_6$ , featuring a hysteresis in the temperature range 313–233 K. Blue and red lines are the line fits for the determination of the activation energies  $E_{A1}$  to  $E_{A3}$ .

units to a tetragonal structure with ordered superoxide  $O_2^-$  and peroxide  $O_2^{2-}$  entities. The  $O_2^-$  units tilt by an angle of  $17^\circ$  relative to the  $c$ -axis in the tetragonal phase, which is a signature of the partially filled molecular orbitals of the paramagnetic  $O_2^-$  units. In contrast to our previous studies, the CO transition was

complete for the present samples of  $Cs_4O_6$ . This was achieved by employing a new synthesis method for the alkali sesquioxides, where the respective superoxides (here  $CsO_2$ ), are thermally decomposed [6].

Due to its molecular nature,  $Cs_4O_6$  has a much lower conductivity than  $Fe_3O_4$ . The electrical conductivity  $\sigma$  was studied by impedance spectroscopy. The structural CO transition is accompanied by an increase in  $\sigma$  of two orders of magnitude, which verifies that a Verwey-type transition occurs. The transition is characterized by a broad hysteresis which is also reflected as an anomaly in the magnetic susceptibility. The various signatures of the Verwey-type transition are represented in Fig. 1. The charge dynamics was investigated on a molecular level by using spectroscopic probes. Particularly useful was  $^{17}O$  NMR spectroscopy which was performed on a sample enriched with the  $^{17}O$  isotope. Due to the large hyperfine fields associated with paramagnetic units, the  $^{17}O$  NMR signal is observed only when diamagnetic  $O_2^{2-}$  ions, which are localized on the NMR time scale, are present. Similarly, an EPR signal is observed only in the presence of localized paramagnetic  $O_2^-$  units. The results obtained from the different experimental techniques are summarized in Fig. 2.

More recently, we could also verify that  $Rb_4O_6$  with smaller  $Rb^+$  ions in its crystal lattice reveals a Verwey-

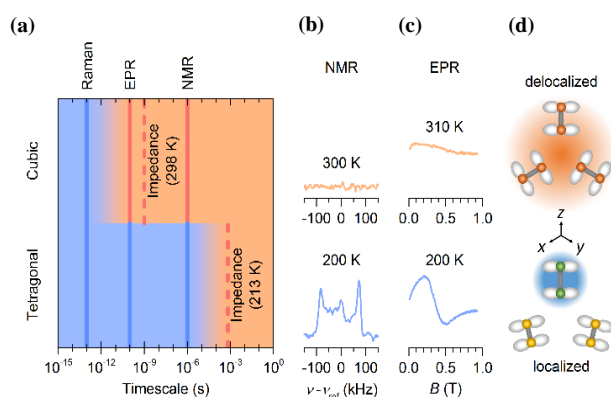


Fig.-2: (from Ref. [5]) (a) Sketch of the characteristic timescales covered by various experimental techniques applied to  $Cs_4O_6$ . (b, c) Due to rapid charge exchange  $^{17}O$  NMR and EPR signals are not observed in the charge-disordered phase (top), whereas both, a  $^{17}O$  NMR signal due to non-magnetic  $O_2^{2-}$  units and an EPR signal due to paramagnetic  $O_2^-$  are apparent in the CO phase (bottom). (d) Schematic illustration of charge disorder with equivalent  $O_2^x$  units and CO with distinctly different  $O_2^{2-}$  (green) and  $O_2^-$  (yellow) units.

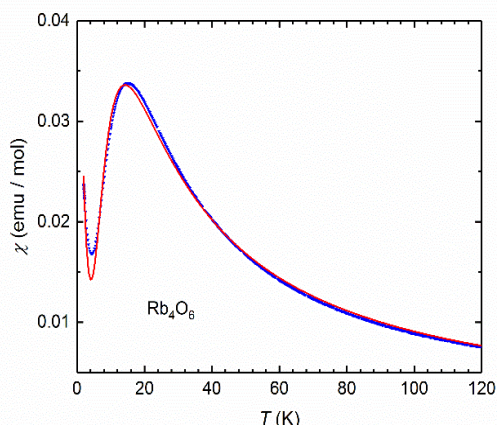


Fig.-3: Temperature dependence of the molar magnetic susceptibility,  $\chi$ , of  $\text{Rb}_4\text{O}_6$  in the low temperature region. The red solid line corresponds to a Bleaney-Bowers fit to the experimental data (blue). The spin Hamiltonian of the dimers is defined as  $H = -2JS_1S_2$  and  $J = -8.5 \text{ cm}^{-1}$  was obtained from the fit. A Curie-term was added to take into account the upturn in  $\chi$  below 5 K.

type CO transition, which is shifted by about 50 K towards higher temperature. In contrast to  $\text{Cs}_4\text{O}_6$ , a hump is seen around 20 K in the low-temperature magnetic susceptibility data (Fig. 3) that can be roughly fitted to the Bleaney-Bowers equation, indicating dimerization of the  $S = \frac{1}{2}$  spins of the  $\text{O}_2^-$  units. The deviations between the experimental and calculated curves suggest that intermolecular interactions play a role as well. Most remarkably, powder neutron diffraction patterns of  $\text{Rb}_4\text{O}_6$  reveal the emergence of new peaks below 90 K, indicating an orbital ordering transition driving the dimerization of  $\text{O}_2^-$  units. Attempts to solve the low-temperature crystal structure, as well as EPR and NMR studies, are in progress in order to understand the relation between orbital ordering and spin arrangements in  $\text{Rb}_4\text{O}_6$ .

The simple CO pattern of molecular  $\text{O}_2^{2-}$  and  $\text{O}_2^-$  units, as well as the lack of long-range magnetic order down to 2 K, suggest the  $\text{A}_4\text{O}_6$  sesquioxides as model systems for disentangling the complex interplay of charge, lattice, orbital, and spin degrees of freedom in Verwey-type CO processes.

### Silver oxoruthenates(V) with low-dimensional structural motifs

Layered transition metal oxides frequently display special electronic instabilities and reveal unexpected magnetic exchange couplings. Here, the two new oxoruthenates(V),  $\text{Ag}_3\text{RuO}_4$  and  $\text{AgRuO}_3$ , with low-dimensional structural motifs were synthesized by

applying hydrothermal methods [7,8]. Their crystal structures (Fig. 4) were determined by single-crystal X-ray diffraction. The crystal structure of  $\beta\text{-Ag}_3\text{RuO}_4$  can be regarded as a ternary, cation-ordered variant of the NiAs type of structure; oxygen replaces As, which consequently adopts a hexagonal close packing (hcp) motif wherein Ag and Ru occupy the octahedral voids in a fully ordered fashion. The cations in  $\beta\text{-Ag}_3\text{RuO}_4$  separate into layers, as can be expressed by the formula  $\text{Ag}(\text{Ag}_{1/2}\text{Ru}_{1/2})\text{O}_2 \equiv \text{Ag}_2[\text{AgRuO}_4]$ . In this manner, the frustration of the antiferromagnetic exchange coupling on the trigonal sublattice hosting Ru(V) has been lifted. The ruthenium(V) ions segregate into clusters of  $(\text{Ru}_4)$  diamonds that are separated by diamagnetic silver ions. The resulting tetrameric polyoxoanions  $[\text{Ru}_4\text{O}_{16}]^{12-}$  are the most prominent local structural constituent. The crystal structure of  $\text{AgRuO}_3$  can be essentially described as a stacking of honeycomb-type  $(\text{RuO}_3)^-$  polyoxoanions along the  $c$ -axis with silver ions located within the van der Waals gap.

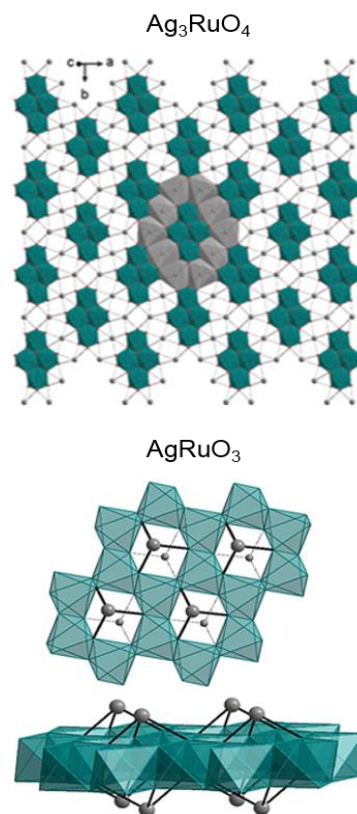


Fig.-4: Illustration of the crystal structures of  $\beta\text{-Ag}_3\text{RuO}_4$  and  $\text{AgRuO}_3$ . The crystal structure of  $\beta\text{-Ag}_3\text{RuO}_4$  features  $[\text{Ru}_4\text{O}_{16}]^{12-}$  tetramers (green) which are separated by silver ions (grey polyhedra). For  $\text{AgRuO}_3$  a single  $(\text{RuO}_3)^-$  slab is shown which builds up the crystal structure [5]. The silver atoms occupy the voids of the structure. The illustrations are adapted from Refs. [7] and [8], respectively.



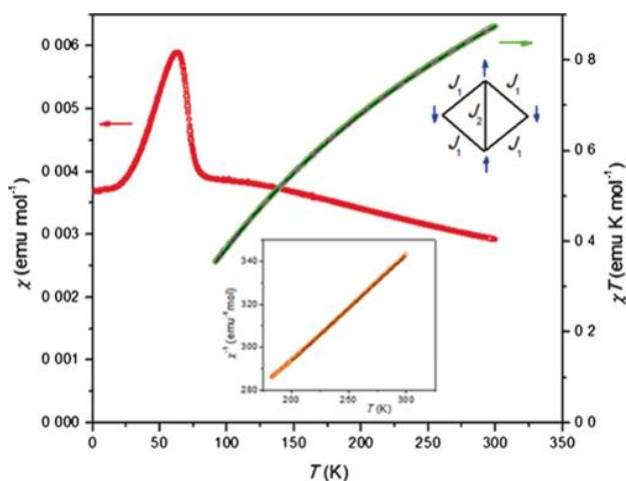


Fig.-5: (from Ref. [7]) Temperature dependence of magnetic susceptibility  $\chi$  (red), reciprocal susceptibility  $\chi^{-1}$  (inset, orange), and product  $\chi T$  (green) of  $\beta$ - $\text{Ag}_3\text{RuO}_4$  under a field of 50 kOe. Fits of  $\chi T$  with the constraint  $J_2=J_1$  (dashed magenta line), and without constraints (black line). Inset, top right: rhomb of ruthenium tetramer, the spin orientation for the ground state is shown.

The magnetic properties of  $\beta$ - $\text{Ag}_3\text{RuO}_4$ ,  $\chi$  and  $\chi T$ , as a function of temperature (Fig. 5), reflect its cluster nature. Above 80 K, within the regime of paramagnetic response,  $\chi$  is reversible and independent of the applied magnetic field. The approximately linear part of  $1/\chi(T)$  between 200 and 300 K follows the Curie-Weiss equation with  $\mu_{\text{eff}} = 4.01 \mu_{\text{B}}$  per Ru atom and  $\theta = -390$  K. This corresponds fairly well to  $\text{Ru}^{5+}$  ions with  $S = 3/2$  ( $\mu_{\text{eff}} = 3.87 \mu_{\text{B}}$ ). It is noteworthy that the large  $\theta$  indicates that strong intra-cluster antiferromagnetic exchange interactions prevail. Around  $\sim 65$  K, a pronounced cusp in the susceptibility suggests a magnetic phase transition. The susceptibility in the range 90–300 K was analyzed using a Heisenberg model with four equal exchange integrals,  $J_1$ , along each side of the rhomb, and different  $J_2$  and  $J_3$  along the short and long diagonals, respectively (see Fig. 5). The ground state of such a tetramer is a spin singlet with two spins pointing down at the long diagonal and two spins pointing up at the short one [7]. Upon cooling below 80 K, after a sharp increase,  $\chi(T)$  exhibits a maximum at about 62–64 K. In AC susceptibility measurements there is no distinct sign of a magnetization relaxation behaviour characteristic of a transition to the spin-glass state. Therefore, it is likely that a transition to a particular magnetically ordered state occurs below 80 K. The field dependence of the magnetization below the transition temperature shows a hysteresis and the coercive field,  $H_c$ , approaches 10 kOe at 1.8 K. The small remanent magnetization,

$M_r$ , which reaches its maximum value of  $0.015 \mu_{\text{B}}$  at 62 K, suggests a canted antiferromagnetic spin structure. The spin canting can be ascribed to anti-symmetric exchange interactions which are a consequence of spin-orbit coupling. The magnetization curves feature a small exchange bias effect, if they are recorded after cooling in a magnetic field. The strong magnetic anisotropy of the  $[\text{Ru}_4\text{O}_{16}]^{12-}$  anion implies that it may reveal the properties of a single-molecule magnet if well-diluted in a diamagnetic matrix.

$\text{AgRuO}_3$  is a semiconductor and, in spite of strong magnetic exchange coupling up to its decomposition temperature of 673 K, neutron diffraction patterns do not show unambiguous evidence of long-range magnetic order. Density functional theory calculations (GGA+U) confirm the 2D character of  $\text{AgRuO}_3$  with large intra-layer exchange coupling,  $J_{ab}$ , and much weaker inter-layer coupling,  $J_c$ .

### External Cooperation Partners

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